

THE GUIDED WAVES MODELLING - A SPECTRAL METHOD APPROACH

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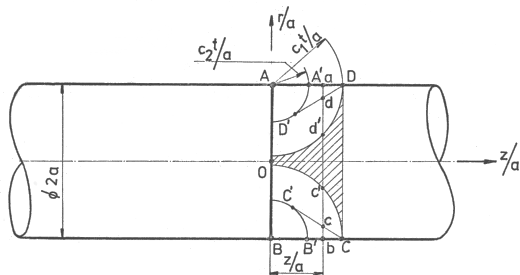
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Motivation



Valeš, F., Podélný ráz polonekonečných válcových elastických tyčí kruhového průřezu, část I (Z681/79) a II (Z839/83), ÚT AVČR Praha, (in Czech).

$$\int_0^{\infty} \frac{(2 - \xi^2) J_1(r \gamma a B) J_1(\gamma a A) - 2 J_1(\gamma a B) J_1(r \gamma a A)}{\gamma a \xi B M} \cos(z \gamma a) \sin(t \xi \gamma a) d\gamma a,$$

where

$$M = \begin{aligned} & -\gamma a^2 ((2 - \xi^2)^2 / A + 4 \kappa A) J_0(\gamma a B) J_0(\gamma a A) \\ & + \gamma a ((2 - \xi^2) (4 + (2 - \xi^2) / (\xi^2 - 1)) + 2 \kappa \xi^2) J_0(\gamma a B) J_1(\gamma a A) \\ & \quad - 2 \gamma a (2 - \xi^2) B / A J_1(\gamma a B) J_0(\gamma a A) \\ & + (2 B (2 \gamma a^2 - (2 - \xi^2) / (\xi^2 - 1)) + \kappa \gamma a^2 (2 - \xi^2)^2 / B) J_1(\gamma a B) J_1(\gamma a A), \end{aligned}$$

a radius of the semi-infinite bar,

γ wavenumber,

ξ the ratio of the phase velocity and the shear wave velocity, $\xi(\gamma a)$,

κ the ratio of the squares of the phase velocities,

$$A \sqrt{\kappa \xi^2 - 1},$$

$$B \sqrt{\xi^2 - 1},$$

J the Bessel function of the first kind.

Characteristic function

The dispersion relations $f(\xi, \gamma a)$ is defined as

$$(2 - \xi^2)^2 J_0(\gamma a A) J_1(\gamma a B) + 4AB J_1(\gamma a A) J_0(\gamma a B) - \frac{2\xi^2}{\gamma a} A J_1(\gamma a A) J_1(\gamma a B) = 0,$$

where

a radius of the semi-infinite bar,

γ wavenumber,

ξ the ratio of the phase velocity and the shear wave velocity,

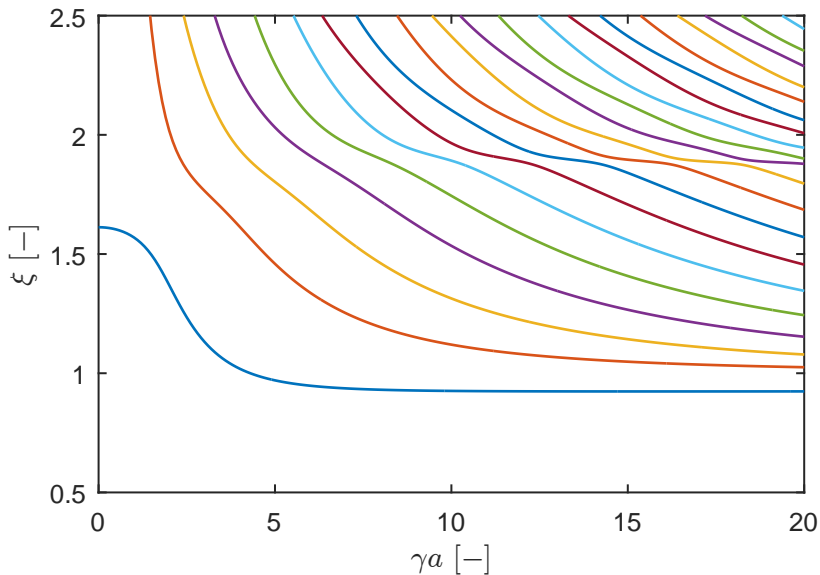
κ the ratio of the squares of the phase velocities,

$$A = \sqrt{\kappa \xi^2 - 1},$$

$$B = \sqrt{\xi^2 - 1},$$

J the Bessel function of the first kind.

Dispersion curves



Solution methods

- ▶ Root-finding

- ▶ Interval arithmetics



Pelikán, Hora:

Robust method for finding of dispersion curves in a thick plate problem, COMPUTATIONAL MECHANICS 2014.

- ▶ Chebyshev interpolation



Hora:

The root-finding of dispersion curves in a bar impact problem, COMPUTATIONAL MECHANICS 2015.



Hora:

The use of the Chebyshev interpolation in elastodynamics problems, COMPUTATIONAL MECHANICS 2016.

- ▶ Spectral method

Equations of motion

Displacement potentials:

$$u_r = \partial_r \phi - \partial_z \psi_\theta,$$

$$u_z = \partial_z \phi + r^{-1} \partial_r (r \psi_\theta).$$

The equations of axisymmetric motion:

$$\nabla^2 \phi = \frac{1}{c_1^2} \partial_t^2 \phi,$$

$$\left(\nabla^2 - \frac{1}{r^2} \right) \psi_\theta = \frac{1}{c_2^2} \partial_t^2 \psi_\theta,$$

where

c_1 the P-wave velocity,

c_2 the S-wave velocity,

∇^2 Laplace operator, $\nabla^2 = \partial_r^2 + r^{-1} \partial_r + \partial_z^2$.

The normal and tangential stress tractions:

$$\begin{aligned}\sigma_{rr} &= \lambda\Delta + 2\mu\partial_r u_r, \\ \sigma_{rz} &= \mu(\partial_z u_r + \partial_r u_z),\end{aligned}$$

where

Δ the dilatation in cylindrical $r - z$ coordinates,
 λ, μ the Lamé parameters.

The propagation of sinusoidal waves:

$$\phi = \Phi e^{i(k_z z - \omega t)}, \quad \psi_\theta = \Psi e^{i(k_z z - \omega t)},$$

where

ω the angular frequency,
 k_z the axial wave number,
 Φ, Ψ the amplitudes.

Helmholtz equations:

$$\underbrace{\left(d_r^2 + r^{-1} d_r + \frac{\omega^2}{c_1^2} \right)}_{\mathcal{L}_1} \Phi = k_z^2 \Phi,$$

$$\underbrace{\left(d_r^2 + r^{-1} d_r - \frac{1}{r^2} + \frac{\omega^2}{c_2^2} \right)}_{\mathcal{L}_2} \Psi = k_z^2 \Psi.$$

Boundary conditions:

$$\sigma_{rr} = \left[-\lambda \left(r^{-2} + \frac{\omega^2}{c_1^2} \right) + 2\mu \partial_r^2 \right] \Phi + 2\mu \partial_r \hat{\Psi},$$

$$\hat{\sigma}_{rz} = -2\mu \left(\partial_r^3 + r^{-1} \partial_r^2 - r^{-2} \partial_r + \frac{\omega^2}{c_1^2} \partial_r \right) \Phi + \mu \left(2\partial_r^2 + 2r^{-1} \partial_r - 2r^{-2} + \frac{\omega^2}{c_2^2} \right) \hat{\Psi},$$

where $\hat{\Psi} = ik_z \Psi$ and $\hat{\sigma}_{rz} = ik_z \sigma_{rz}$.

Spectral method

- ▶ Differential equations can be very efficiently solved with spectral collocation methods. Orthogonal polynomials of high degree are used as **global** interpolants to approximate the unknown functions of considered differential equations.
- ▶ The discrete matrix operator, which approximates the differential operator, is called a **differentiation matrix**.
- ▶ The differentiation matrix can be based on **Chebyshev**, Fourier, Hermitian, or other interpolants, which can be differentiated exactly.
- ▶ The global interpolant evaluated at N interpolation points is connected to its first derivative by a matrix vector product.
- ▶ Evenly spaced points \rightarrow the Runge phenomenon.
- ▶ Chebyshev points:

$$x_i = \cos\left(\frac{(i-1)\pi}{N-1}\right), \quad i = 1, \dots, N$$

$$f(x) \approx \sum_{k=1}^{N+1} f(x_k) T_k(x)$$

$$f(x_j)^{(\ell)} = \sum_{k=1}^{N+1} \frac{d^\ell}{dx^\ell} (T_k(x_j)) f(x_k)$$

$$D_{jk}^{(\ell)} = \frac{d^\ell}{dx^\ell} (T_k(x_j))$$

⇓

$$\begin{pmatrix} f_1^{(m)} \\ f_2^{(m)} \\ \vdots \\ f_N^{(m)} \end{pmatrix} \approx \begin{pmatrix} D_{11}^{(m)} & D_{12}^{(m)} & \cdots & D_{1N}^{(m)} \\ D_{21}^{(m)} & \ddots & \ddots & \vdots \\ \vdots & \ddots & \ddots & \vdots \\ D_{N1}^{(m)} & \cdots & \cdots & D_{NN}^{(m)} \end{pmatrix} \cdot \begin{pmatrix} f_1 \\ f_2 \\ \vdots \\ f_N \end{pmatrix}$$

Chebyshev differential matrix, resources

- ▶ CHEBFUN (MATLAB toolbox)



Driscoll, Hale, Trefethen, editors, Chebfun Guide, Pafnuty Publications, Oxford, 2014.

- ▶ CHEBDIF (MATLAB routine)



Weideman, Reddy, A MATLAB differentiation matrix suite, ACM Trans. Math. Softw. 26, 2000.

- ▶ ApproxFun (Julia package)



URL: github.com/ApproxFun/ApproxFun.jl.git

- ▶ pychebfun (Python implementation of chebfun)



URL: github.com/oliviervedier/pychebfun

Matrix forms

Helmholtz equations:

$$L_1 = D^{(2)} + \text{diag} \left(\frac{1}{r} \right) D^{(1)} + \left(\frac{\omega^2}{c_1^2} \right) \mathbf{I}$$

$$L_2 = D^{(2)} + \text{diag} \left(\frac{1}{r} \right) D^{(1)} - \left(\frac{1}{r^2} \right) \mathbf{I} + \left(\frac{\omega^2}{c_2^2} \right) \mathbf{I}$$

The stress components:

$$S_{r\Phi} = -\lambda \left[\text{diag} \left(\frac{1}{r^2} \right) + \left(\frac{\omega^2}{c_1^2} \right) \mathbf{I} \right] + 2\mu D^{(2)},$$

$$S_{r\hat{\Psi}} = 2\mu D^{(1)},$$

$$S_{z\Phi} = -2\mu \left[D^{(3)} + \text{diag} \left(\frac{1}{r} \right) D^{(2)} - \text{diag} \left(\frac{1}{r^2} \right) D^{(1)} + \left(\frac{\omega^2}{c_1^2} \right) \mathbf{I} \right],$$

$$S_{z\hat{\Psi}} = \mu \left[2D^{(2)} + 2\text{diag} \left(\frac{1}{r} \right) D^{(1)} - \text{diag} \left(\frac{1}{r^2} \right) + \left(\frac{\omega^2}{c_2^2} \right) \mathbf{I} \right].$$

Combination of equations

Helmholtz equations:

$$\underbrace{\begin{pmatrix} L_1 & 0 \\ 0 & L_2 \end{pmatrix}}_L \begin{pmatrix} \Phi \\ \hat{\Psi} \end{pmatrix} = k_z^2 \begin{pmatrix} \Phi \\ \hat{\Psi} \end{pmatrix}$$

The stress components:

$$\begin{pmatrix} \sigma_{rr} \\ \hat{\sigma}_{rz} \end{pmatrix} = \underbrace{\begin{pmatrix} S_{r\Phi} & S_{r\hat{\Psi}} \\ S_{z\Phi} & S_{z\hat{\Psi}} \end{pmatrix}}_S \begin{pmatrix} \Phi \\ \hat{\Psi} \end{pmatrix}$$

Embedding boundary conditions

1. The lines in the L matrix corresponding to $r = a$ will be replaced by the corresponding lines of the S matrix.
2. The corresponding values on the right-hand have to be set equal to zero.

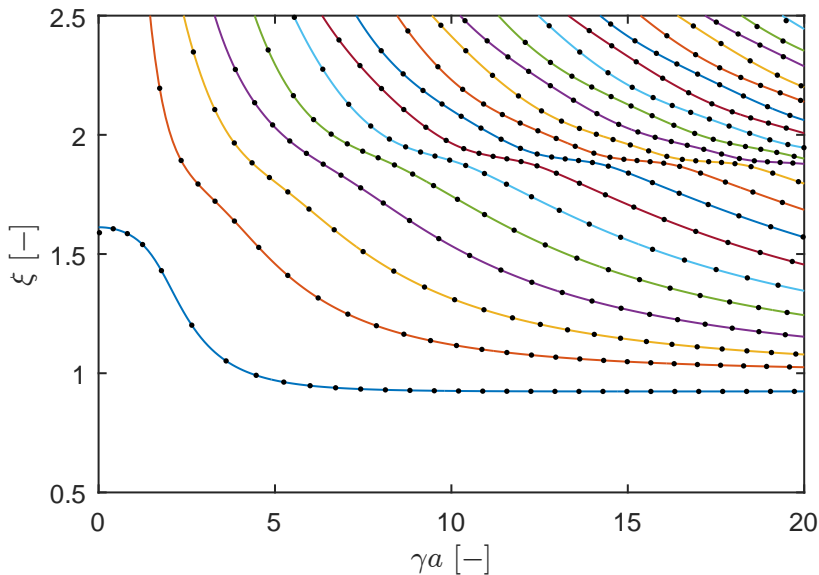
$$\begin{pmatrix} L_1 & 0 \\ 0 & L_2 \end{pmatrix} \begin{pmatrix} \Phi \\ \hat{\Psi} \end{pmatrix} = k_z^2 \begin{pmatrix} M & 0 \\ 0 & M \end{pmatrix} \begin{pmatrix} \Phi \\ \hat{\Psi} \end{pmatrix},$$

where

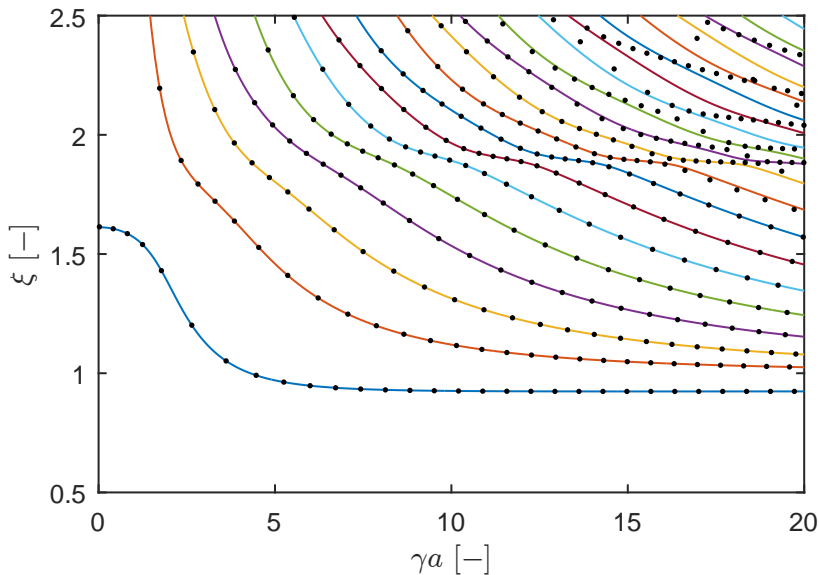
$$M = \begin{pmatrix} 0 & & & & \\ & 1 & & & \\ & & \ddots & & \\ & & & 1 & \\ & & & & 0 \end{pmatrix}$$

Generalized eigenvalue problem \rightarrow MATLAB routine `eig`.

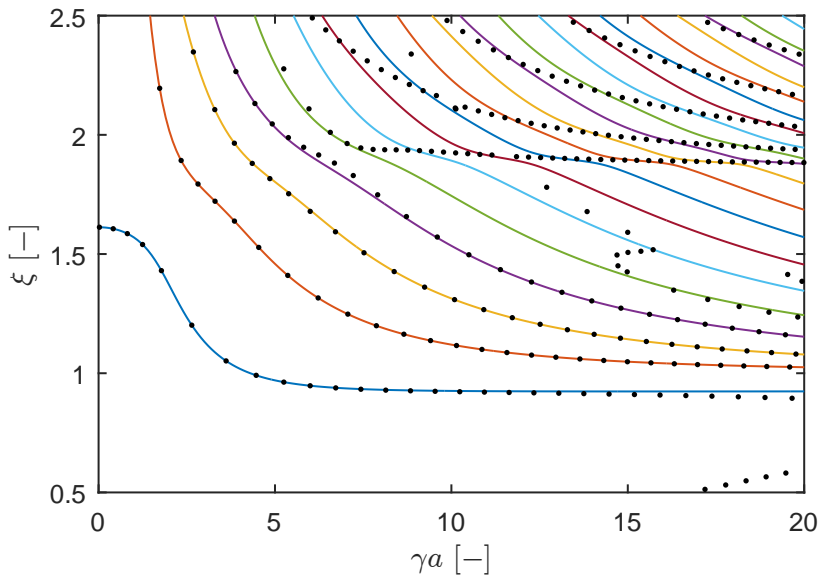
Dispersion curves - spectral method, $N = 30$



Dispersion curves - spectral method, $N = 20$



Dispersion curves - spectral method, $N = 10$



Thank you for your attention!

Any questions?

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